Developing a nonhydrostatic isopycnalcoordinate ocean model

Oliver Fringer

Associate Professor

The Bob and Norma Street Environmental Fluid Mechanics Laboratory Dept. of Civil and Environmental Engineering

Stanford University

11 September 2015 Funding: ONR Grants N00014-10-1-0521 and N00014-08-1-0904

Outline

- Overview of internal gravity waves
- Nonhydrostatic (Navier-Stokes) modeling
- Grid-resolution requirements
 →Nonhydrostatic modeling is expensive!
- A nonhydrostatic isopycnal-coordinate model
 →The cost can be reduced!
- Conclusions

Internal gravity waves



- Typical speeds in the ocean: 1-3 m/s
- Frequencies: Tidal (internal tides) minutes (internal waves)
- Wavelengths: 100s of km to 10s of m



Applications of internal gravity waves

 Breaking of internal tides and waves may provide the necessary mixing to maintain the ocean stratification (Munk and Wunsch 1998).
 warm equatorial



Mixing induced by breaking internal waves prevents the ocean from turning into a "stagnant pool of cold, salty water"...

- Internal waves are hypothesized to deliver nutrients that sustain thriving coral reef ecosystems (e.g. Florida Shelf, Leichter et al., 2003; Dongsha Atol, Wang et al. 2007)
- Internal waves influence sediment transport in lakes and oceans and propagation of acoustic signals.
- Strong internal wave-induced currents can cause oil platform instability and pipeline rupture.

Isopycnal vs z- or sigma-coordinates



isopycnal coordinates

Advantages of isopycnal coordinates:

• Reduces the number of vertical grid points from O(100) in traditional coordinates to O(1-10)

• No spurious vertical (diapycnal) diffusion/mixing

Challenges of isopycnal coordinates:

- Cannot represent unstable stratification
- Layer outcropping (drying of layers) requires special numerical schemes
- Hydrostatic

Hydrostatic vs. nonhydrostatic flows

- Most ocean flows are hydrostatic
 - Long horizontal length scales relative to vertical length scales, i.e. long waves (i.e. Lh >> Lv)
- Only in small regions is the flow nonhydrostatic
 - Short horizontal length scales relative to vertical scales (i.e. Lh ~ Lv)
 - Can cost 10X more to compute!



Nonhydrostatic effects: Overturning

Hydrostatic



Nonhydrostatic





Overturning motions and eddies are not the only nonhydrostatic process...

Nonhydrostatic effects: Frequency dispersion of gravity waves

• Dispersion relation for irrotational surface gravity waves:

$$c^2 = \frac{g}{k} \tanh kD = \frac{g}{k} \tanh \pi \varepsilon, \quad \varepsilon = \frac{D}{L}$$

Deep-water limit: ε>>1 (nonhydrostatic)

$$c^2 = \frac{g}{k}$$

Shallow-water limit ε<<1 (hydrostatic)

$$c^2 = gD$$



When is a flow nonhydrostatic?

Nonhydrostatic Model

Hydrostatic Model

11



Aspect Ratio:

 $\varepsilon = D/L = 2$

,		l		1	ļ	l	1	1	1		١.	1	1	1	1		1	1							2		-																															
		Ì	l	1	j		1	1	1		ľ	¥	1	1		ł	4	3	3	1))	0	¥.	1	>	~	•	-	-			/			1	1	1	1	1		r	t	t	İ	1		4	4	¢	4	1	И	V	V	И	Ņ	1
		ļ	j		l	ļ	1	l		l	۹ ١	۴ ۱	1	1	1	1	4	4	4	0	1	۱	0	<i>.</i>	`	`	i N	•	•	-	4	•	/		,	1	1	1	1	1	1	ŕ	t	f	1	Ì	ì	h	f	ſ	ĵ	j	1	ų	V	M	N	1
		I	ļ		l	ļ	Ì	ł	i i		ľ	ľ	1	1	1		1	4	3		•	2	1		`	`	0	•	•	-			/		•	ſ	1	1	1	1	1	ŕ	1	f	1	1	1	þ.	ŕ	ſ	1	1	ł	ų	1	ł	2	l
		1	1	I	I	l	ļ	ł	1		۹ ۱	۹ ۱	1	1	1		4	4	4	0 	•)		1.	`	1	1		•	-		-	/		•	/	1	/	1	1	1	ŕ	1	ſ	1	1	1	h	ŕ	ſ	î	1	ł	ų	1	l	.,	l
		ļ	J	l	l	ļ	ļ	ł	.)	e R	۹ ١	ľ.	s A	1	i N		1	1	3	1 	• •)	e 	¥.		`		•	-	-			/		•	/	1	l	1	1	1	ŕ	1	f	1	1	ì	þ.,	t	î	f	1	1	l	1	ł	ľ	i
		ļ	J	1	j	1	J	Î	.)		۹ ١	۴	1	1	1		1	•	1	e S	•	2	0		`	1		**	•	-		-	/		•	/	1	ł	1	1	1	f	t	f	1	1	ĩ	A.	t	î	t	1	ľ	ų	1		L	I
		J	J	1	1	ļ	ļ	1	1	l	ľ	4	1	1			•	1	2	•	•)			`	`	`	•	•	-		~	/	2	•	/	1	f	1	1	1	t	t	f	1	1	ĩ	-	t	î	t	1	ľ	1	1		ľ	l
		J	1	1	1	ļ	1	1	.1	1	ĺ	1	1				1		2		1	2			`	'	-		•	-		1	/		•	,	1	1	7	7	1	t	t	t	t	1	ĩ	1	t	î	ţ	1	1	1	1	ł	1	I
		J	1		1	1	J	1	.)		ì	1	1	1	6 1	8			1) 	2	2		×	>				•	-		•			•	,	1	1	1	7	1	f	1	f	ł	1	1	į	ť	î	f	1	ľ	1	í	l	Ì	l
		ļ	4	1	i	1	ļ	1			1	ì	1	1					1		1				2				*	-					•	\$	1	/	1	'	1	1	t	f	ł	1	1	1	ŕ	t	t	1	1	1	1	1	•	l
		1	1	1	1	1	1	1	1		ì	1	1	. 1	ii N		1									2	~		•			-	/	1		'	,	1	1	';	1	t.	1	t	t	1	1		t	t	t	1	1	1	ï	1		l
		1	4	1	Ļ	1	1	1	4	i	í	1	١				ŝ	1	Ň		î			8	2 0					-				2	•	/	,	/	1	')	1	1	1	1	t	1	1	1	t	t	t	1	ĺ	1	ï	ł	1	l
		1	4	1	Ļ	1	4	1	d		1	Ň	1	1		1	1	1	1	0	Ì				2	2					Ĭ					/	/	/	1	'	1	/	1	1	1	1	1		t	1	t	1	Ì	1	1		1	i
		1	4	1	4	4	4	.1	1	Ì	1	1	1	1				1	1		×																	1	1	2			ť	2	1	1	1		t	1	t	1	ł	1	1	ł	1	i
		ļ	4	1	\$	1	4	.1	h	Ŕ	4	Ń	1			5	ŝ	~				,																				() 	1		1	1	1		!	1	t.	1		1	1	l		i
	1	+	4	1	4	4	1	1	ň	í	Ň	1	1		12		8	~				,														,	2					2	,	i G	1	1	1		8 2	1		1	5	1		l		1
		1	4	1	ş	\$	4	,	1	i	,		1		ñ		~	/			7									_		_										,		,	,	í	7		r r	7	,	7	1	1		8		
2	1	1	3	1	3	4		1		į	•	`	`					,	,		,	,							-	_			,			,	,	,	.,	.,		,	,	,	: {	1			ĩ	1	1	1		1				
		4	3	4	,	4	1	,	0	ł	•	1	`				ş	`		•	,	,		•	~	2	ł					 •				-	~	,		.,			,	,	,	3	1		,	,	,	,		2 9				,
	i.	1	5	4	6	3	,	,			•	`	,				,	,		• •	,	1			-			2	•						-		-	,				6		,	,				,	,	,	,	6			ě.	l	,
		1	-1	•	4	,	1			ŝ		`	,			1	2	,			,					,											-			.,		1		,	,	,			è	,	,	,		4				,
		,	•		•	,				3		~				5		-	~			-				-			-			-	-				-	-				-	-		,				•	,	,	,		,				•
																-	-	-	-			-		-	-	-			-	-	-	-	-		-		-					2	-	-	-		-		-	-					1	-		
								ĺ												ľ	ſ																																		1		1	



Aspect Ratio:

$$\varepsilon = D/L = 1/8 = 0.125$$

Nonhydrostatic result = Hydrostatic result + ε^2

Example 3D nonhydrostatic z-level simulation: Internal gravity waves in the South China Sea



Generation of weakly nonlinear wavetrains



Long internal tides $O(100 \text{ km}) \rightarrow \text{Short}$, solitary-like waves O(5 km)

How can we determine, apriori, how much horizontal grid resolution is needed to simulate this process?

Internal solitary waves



Nonlinear effect (steepening): $\delta = a/h_1$ Nonhydrostatic effect (frequency dispersion): $\epsilon = h_1/L$

Solitary wave:
Balance between nonlinear steepening
and nonhydrostatic dispersion. $\delta \sim \epsilon^2$

The KdV equation

When computing solitary waves, the behavior of a 3D, fully nonhydrostatic ocean model can be approximated very well with the KdV (Korteweg and de-Vries, 1895) equation:



The KdV equation gives the well-known solution

Numerical discretization of KdV

- Many ocean models discretize the equations with second-order accuracy in time and space.
 (e.g. SUNTANS, Fringer et al. 2006; POM, Blumberg and Mellor, 1987; MICOM, Bleck et al., 1992; MOM, Pacanowski and Griffes, 1999).
- A second-order accurate discretization of the KdV equation using leapfrog (i.e. POM) is given by

$$\frac{\partial\xi}{\partial t} + \left(1 - \frac{3}{2}\delta\xi\right)\frac{\partial\xi}{\partial x} + \frac{\epsilon^2}{6}\frac{\partial^3\xi}{\partial x^3} = 0$$

$$\frac{\xi_i^{n+1} - \xi_i^{n-1}}{2\Delta t} + \left(1 - \frac{3}{2}\delta\xi_i^n\right)\frac{\xi_{i+1}^n - \xi_{i-1}^n}{2\Delta x} + \frac{\epsilon^2}{6}\frac{\frac{1}{2}\xi_{i+2}^n - \xi_{i+1}^n + \xi_{i-1}^n - \frac{1}{2}\xi_{i-2}^n}{(\Delta x)^3} = 0$$

• Use the Taylor series expansion to determine the modified equivalent form of the terms, e.g.

$$\frac{\xi_{i+1}^n - \xi_{i-1}^n}{2\Delta x} = \frac{\partial \xi}{\partial x} \Big|_i^n + \frac{(\Delta x)^2}{6} \frac{\partial^3 \xi}{\partial x^3} \Big|_i^n + \frac{(\Delta x)^4}{120} \frac{\partial^5 \xi}{\partial x^5} \Big|_i^n + \frac{(\Delta x)^6}{5040} \frac{\partial^7 \xi}{\partial x^7} \Big|_i^n + O\left((\Delta x)^8\right)$$

Modified equivalent KdV equation

The discrete form of the KdV equation produces a solution to the modified equivalent PDE (Hirt 1968) which introduces new terms due to discretization errors

KdV
$$\frac{\partial \xi}{\partial t} + \left(1 - \frac{3}{2}\delta\xi\right)\frac{\partial \xi}{\partial x} + \frac{\epsilon^2}{6}\frac{\partial^3 \xi}{\partial x^3} = 0$$

ent KdV:

Modified equivalent KdV:

$$\frac{\partial\xi}{\partial t} + \left(1 - \frac{3}{2}\delta\xi\right)\frac{\partial\xi}{\partial x} + (1 + \Gamma)\frac{\epsilon^2}{6}\frac{\partial^3\xi}{\partial x^3} = O\left(\epsilon^2(\Delta x)^2, \delta(\Delta x)^2, (\Delta x)^4, (\Delta t)^4\right)$$
$$\Gamma = K\left(\frac{\Delta x}{h_1}\right)^2 = \frac{\text{Numerical dispersion}}{\text{Physical dispersion}} = K\lambda^2 \qquad \begin{array}{c} \lambda = \Delta x/h_1 \text{ grid "lepticity"} \\ \text{(Scotti and Mitran, 2008)} \\ \text{K} = O(1) \text{ constant.} \end{array}$$

The numerical discretization of the first-order derivative produces numerical dispersion. Note that the errors in the nonlinear term are smaller by a factor δ .

For numerical dispersion to be smaller than physical dispersion, $\lambda < 1$.

Hydrostatic vs. nonhydrostatic for λ =0.25



Numerical dispersion is 16 times smaller than physical dispersion.

Vitousek and Fringer (2011)

Hydrostatic vs. nonhydrostatic for λ =8



Vitousek and Fringer (2011)

Nonhydrostatic isopycnal model?

- Zhang et al. simulation: 12 million cells, ∆x=1 km = 5 h₁
- To *begin* to resolve nonhydrostatic effects, ∆x=200 m = h₁ → 300 million cells! With ∆x=100 m, 1.2 billion!
- The z-level SUNTANS model requires O(100) z-levels to minimize numerical diffusion of the pycnocline.
- Solution: Isopycnal model with O(2) layers
 = 50X reduction in computation time.
- → Nonhydrostatic isopycnal coordinate model.

2-layer hydrostatic result with isopycnal model of Simmons, U. Alaska Fairbanks.



Essential features of the nonhydrostatic isopycnal-coordinate model



- Staggered C-grid layout
- Split Montgomery potential into Barotropic (implicit) & Baroclinic (explicit) parts
- IMplicit-EXplicit (IMEX) multistep method. Durran (2012)
- MPDATA for upwinding of layer heights
- Implicit theta method for vertical diffusion
- Explicit horizontal diffusion
- Predictor/corrector method for nonhydrostatic pressure
- \rightarrow Second-order accurate in time and space

Vitousek and Fringer (2014)

Nonhydrostatic test cases



Note: density need not change in each layer

Hydrostatic internal seiche



2 layers

100 layers

Nonhydrostatic internal seiche

Horizontal Velocity [m/2]



2 layers

100 layers

Dispersion relation speed = function(wavelength)



Internal solitary wave formation



z-level model leads to numerical diffusion, or thickening of the pycnocline.

Isopycnalcoordinate model:

- → eliminates spurious numerical diffusion
- → captures solitary wave behavior at 1/50 cost...

Vitousek and Fringer (2014)



Internal Solitary Waves in the (idealized) South China Sea

Vitousek and Fringer (2014)

Internal Wave Runup



Vitousek and Fringer (2014)

Conclusions

- Simulation of nonhydrostatic effects in the SCS requires $\Delta x < h_1 \rightarrow O(billion)$ grid cells in 3D with z-level model.
- We have developed a nonhydrostatic isopycnalcoordinate model using stable higher-order timestepping.
- More isopycnal layers are needed:
 - To resolve stratification
 - To resolve nonhydrostatic effects
- Most oceanic/lake processes are weakly nonhydrostatic and so <10 layers suffice for many applications. The result is a reduced computational cost by O(10).
- Ongoing work: Development of unstructured-grid model.